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Spectral analysis of the stochastic Helmholtz equation in the plane

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1 Introduction

We consider time-harmonic acoustic or transverseelectric (TE) electromagnetic waves in the plane at frequency $k_0 > 0$. A source f is supported in a subset of an open, relatively compact set $\Omega \subset \mathbf{R}^2$ and radiates a wave *u* in the presence of a stochastic medium. The wave is governed by an inhomogeneous Helmholtz equation and the Sommerfeld radiation condition,

 $\begin{cases} \left(\Delta + (1+q(x))k_0^2\right)u = f & \text{in } \mathbf{R}^2, \\ \lim_{|x| \to \infty} \sqrt{|x|}(\partial_{|x|}u - ik_0u) = 0 \end{cases}$ uniformly in $x/|x| \in S^1$.

The contrast q is a second-order random field with each realization a.s. in $L^2(\Omega)$. We fix p > 2 and require $f \in L^p(\Omega)$, ensuring [1] the existence and uniqueness of solution of (*) in $L^p(\Omega)$. Furthermore, the contrast satisfies

The operator V_1 has a known singular value decomposition [2, 3], and we here consider the normal operator for $V_a V_1$,

 $\Upsilon = (V_q V_1)^* (V_q V_1).$

Iterated application of the expansion of the fundamental solution Φ of the Helmholtz operator,

$$\Phi(x-y) = \sum_{n \in \mathbb{Z}} H_n^{(1)}(\kappa) J_n(k_0|y|) e^{in(\arg(x) - \arg(y))}$$

yields [4] a decomposition of Υ to

$$\begin{split} \Upsilon f &= \sum_{n \in \mathbb{Z}} \tilde{\lambda}_n \langle f, \psi_n \rangle_{L^2(\Omega)} \psi_n + \\ &\sum_{n, m \in \mathbb{Z}} \lambda_{nm} \langle f, \psi_n \rangle_{L^2(\Omega)} \psi_m \end{split}$$

with 'spectral values' $\tilde{\lambda}_n$, λ_{nm} , and functions ψ_m , $\psi_m(\cdot) = J_m(k_0|\cdot|)e^{im \arg(\cdot)},$ $\lambda_{nm} = (2\pi R)^2 \,\sigma H_n^{(1)}(\kappa) \xi_m \Big| H_m^{(1)}(\kappa) \Big|^2 H_m^{(2)}(\kappa) J_m(\kappa),$ $\tilde{\lambda}_n = (2\pi R)^2 \eta \beta_n(\kappa) |H_n^{(1)}(\kappa)|^4 J_n(\kappa).$



1. $\operatorname{supp}(f) \subseteq \operatorname{supp}(q) \subset \Omega$ 2. q(x) > 0 a.s, $q|_{\partial\Omega} = \eta = \text{const.}$ In particular,

 $q = \eta + \sigma \frac{dW(x)}{dx} \mathbb{1}_{\Omega}$

where dW(x)/dx is the formal derivative of the two-sided Brownian sheet. We are here interested in the transfer of information from f to u, in terms of the singular spectrum of the mapping $f \mapsto u$.

2 Forward and normal operator

The system (*) is equivalent with the Lippmann-Schwinger equation

 $\left(I - k_0^2 V_q\right)u = V_1 f$

with the integral operator V_q

$$(V_q f)(x) = \int_{\Omega} \Phi(x - y) q(y) f(y) dy.$$

For wavenumber k_0 fulfilling

Here

$$\mathcal{B}_n(\kappa) = \int_0^R r J_m(k_0 r) dr \delta_{n,m}$$

and ξ_m are normally distributed $\mathcal{N}(0, ||\psi_m||^2_{L^2(\Omega)})$. In the second term of the decomposition of Υ , we discover 'spectral leakage.'

3 Numerical results

With $k_0 = 2\pi$ and $\Omega = \{x \in \mathbb{R}^2, |x| < 5\}$, the values $\eta = \sigma = 0.001$ ensure the validity of (1). In this case the values $|\tilde{\lambda}_n|$ are as shown below.



This does not yield a spectral decomposition of the operator V_qV_1 . However, for any given source f it is possible to numerically estimate the singular values of V_qV_1 utilizing the normal equations.

4 Conclusion

The source-to-field measurement operator for the inhomogeneous Helmholtz equation with random wavenumber has a singular spectrum with a nontrivial random component inducing spectral leakage. This not only influences the magnitude of the spectral values but also causes a nonuniform spatial frequency phase-shift in the measurement. These effects combine to make the inverse source/inverse medium problem even more ill-posed in the presence of a random medium.

Bibliography

[1] Lechleiter, Kazimierski and Karamehmedović, Inverse Probl. 29 (2013) 075003

$$k_0 < \left(||q||_{L^{\infty}(\Omega)} ||\Phi(x-\cdot)||_{L^2(\Omega)} \right)^{-1/2}, \quad (1)$$

the inverse of $I - k_0^2 V_q$ is expressible in terms of a convergent Neumann series. Truncating the series at first order defines the mapping $f \mapsto u$ as

 $u = (V_1 + k_0^2 V_q V_1) f + O(||q||_{L^{\infty}(\Omega)}^2).$

Furthermore, here are the mean magnitude and mean complex argument, respectively, of 400 realizations of λ_{nm} :

[2] Karamehmedović, J. Phys. Commun. 2 (2018) 095021

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